

# A Theorist's Radio Sky

*Xuelei Chen*

@NAOC

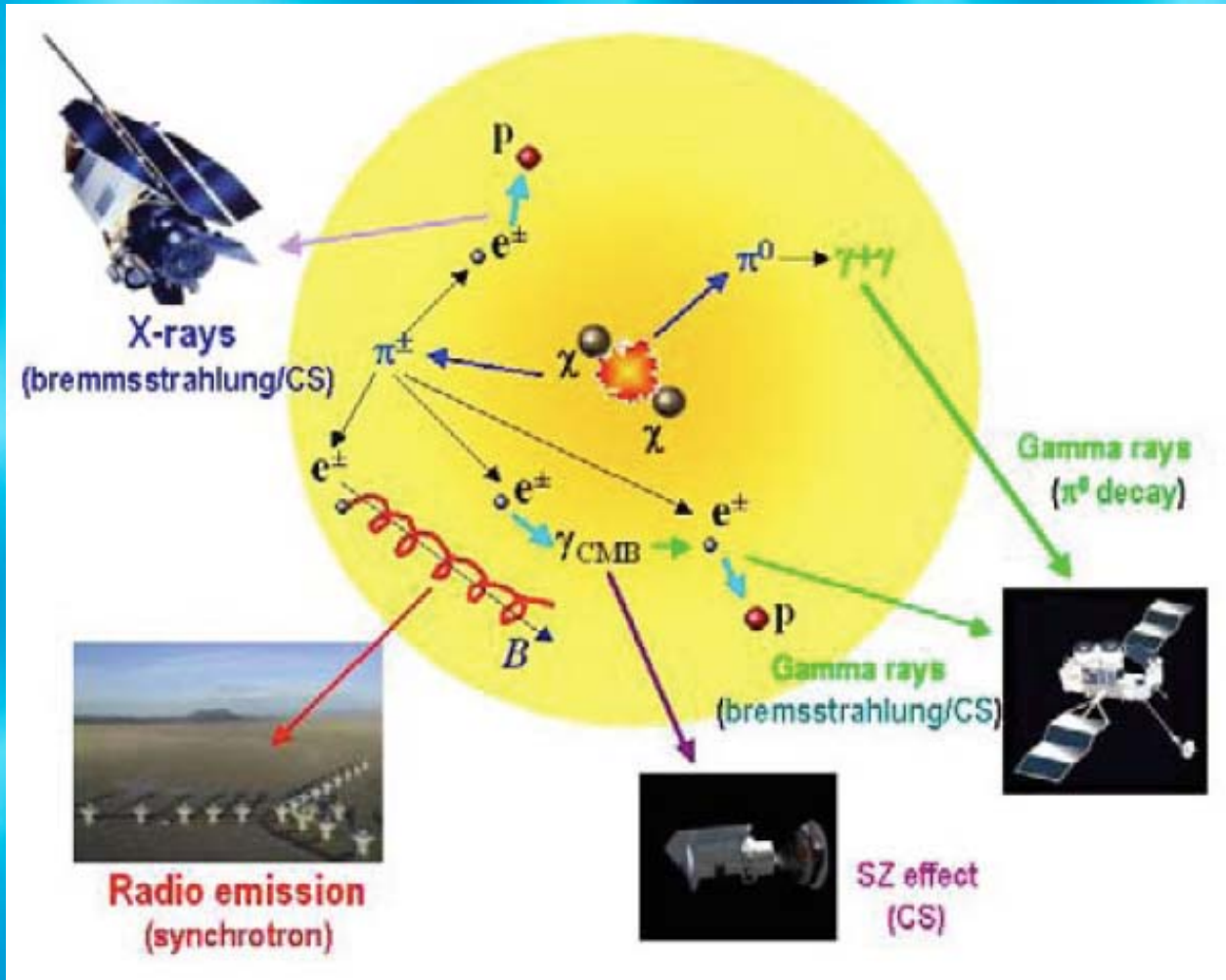
# Outline

radio detection of dark matter

21cm signature of first stars

21cm forest

# DM annihilation signals



## Electron & Positron Diffusion in magnetic field

$$n_e(E, r) = [Q_e(E, r)\tau_{loss}] \cdot \frac{V_{source}}{V_{source} + V_{diffusion}} \cdot \frac{\tau_D}{\tau_D + \tau_{loss}}$$

$$\tau_{loss} \gg \tau_D$$

$$n_e(E, r) = [Q_e(E, r)\tau_{loss}] \cdot \frac{V_{source}}{V_{diffusion}} \cdot \frac{\tau_D}{\tau_{loss}}$$

# Possible local DM annihilation sources

**Annihilation rate**  $\propto \rho^2$

## ⊙ Galactic Center

nearest and densest region

but with numerous astrophysical activities

## ⊙ Nearby galaxies cluster: Coma

cluster as the largest bound containers of DM

## ⊙ Dwarf galaxies: dwarf spheroidal galaxies

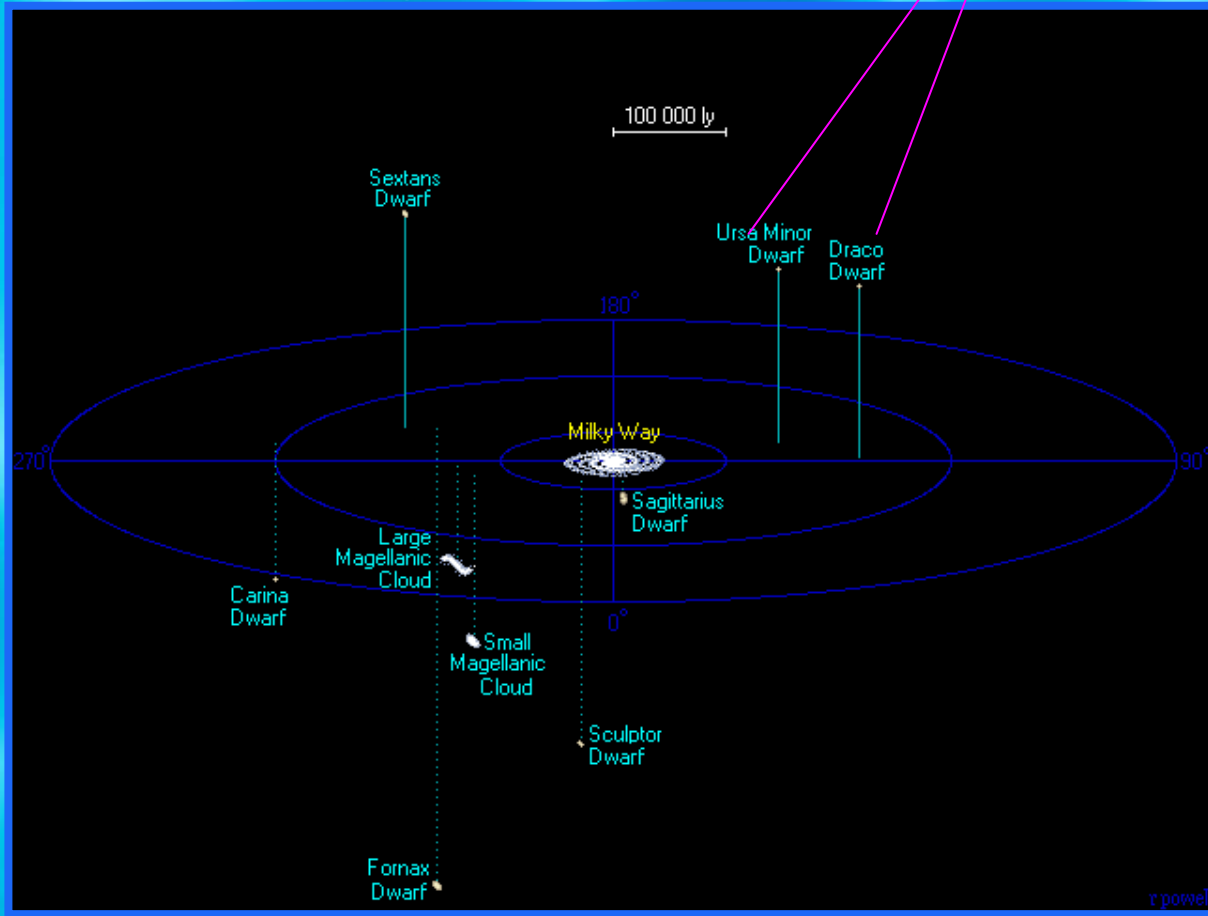
dominated by DM both in central regions and outskirts;

close to be a *pure* dark halo;

local group satellites

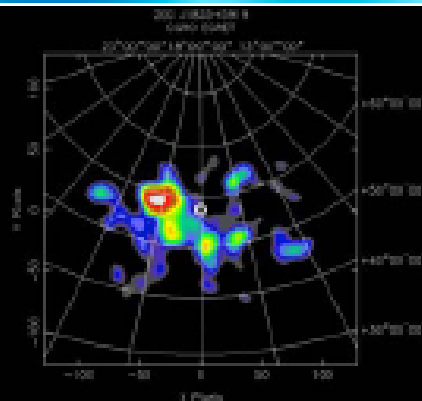
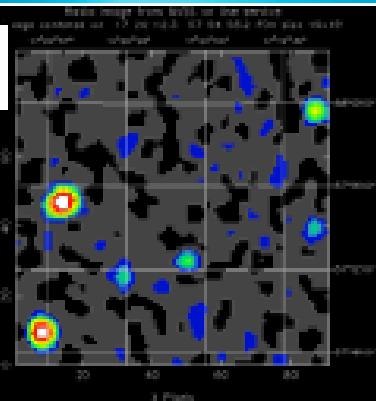
# Milk Way dSphs

High latitude  
Highest mass/light



# Draco and Ursa Minor

Draco

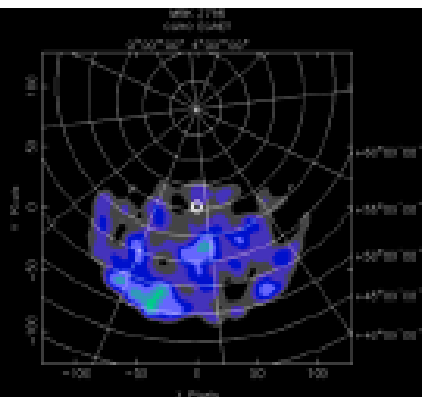
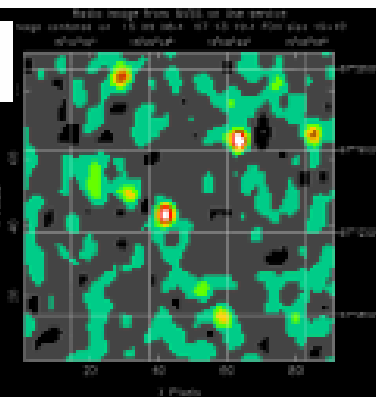


$$F_{4.9} < 2 \text{ mJy}$$

$$F_{2\text{keV}} < 1.7 \times 10^{-14} \text{ erg/cm}^2$$

$$S=1 \text{ pho; bkg}=1 \text{ pho}$$

U Mi

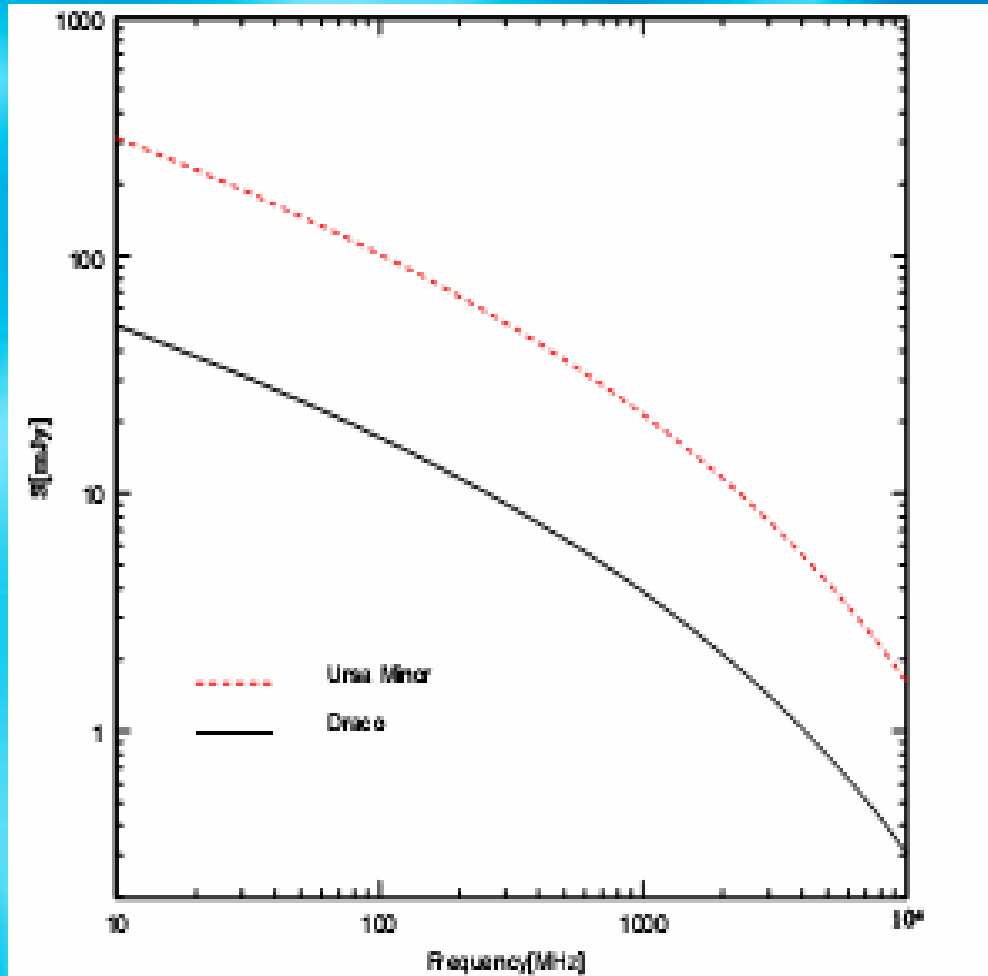


$$F_{1.4} < 3.5 \text{ mJy}$$

$$F_{2\text{keV}} <$$

$$S=13 \text{ pho; bkg}=11 \text{ pho}$$

# Estimation of the radio emission:



## 1) Dark halo:

$$\rho_{NFW}(r) = \frac{\rho_s}{(r/r_s)(1+r/r_s)^2}$$

with two parameters determined from the velocity dispersion of stars

## 2) Magnetic field:

$$1 \mu\text{Gs}$$

## 3) neutralino:

$$M_\chi = 100\text{GeV}$$

$$\langle \sigma v \rangle \simeq 2.0 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$



# Diffuse radio emission in dSphs

Table 2. Related results ( $\theta$ : half of the angular diameter)

|                | Flux(mJy) from Draco |         |        | Flux(mJy) from Ursa Minor |         |        |
|----------------|----------------------|---------|--------|---------------------------|---------|--------|
|                | 4.89GHz              | 1.42GHz | 0.7GHz | 4.89GHz                   | 1.42GHz | 0.7GHz |
| $\theta = 6'$  | 0.2                  | 0.7     | 1.1    | 0.4                       | 1.5     | 2.7    |
| $\theta = 30'$ | 0.8                  | 2.9     | 5.1    | 4.4                       | 15.9    | 28.6   |
| $\theta = 60'$ | 0.9                  | 3.2     | 5.7    | 6.2                       | 22.8    | 40.9   |

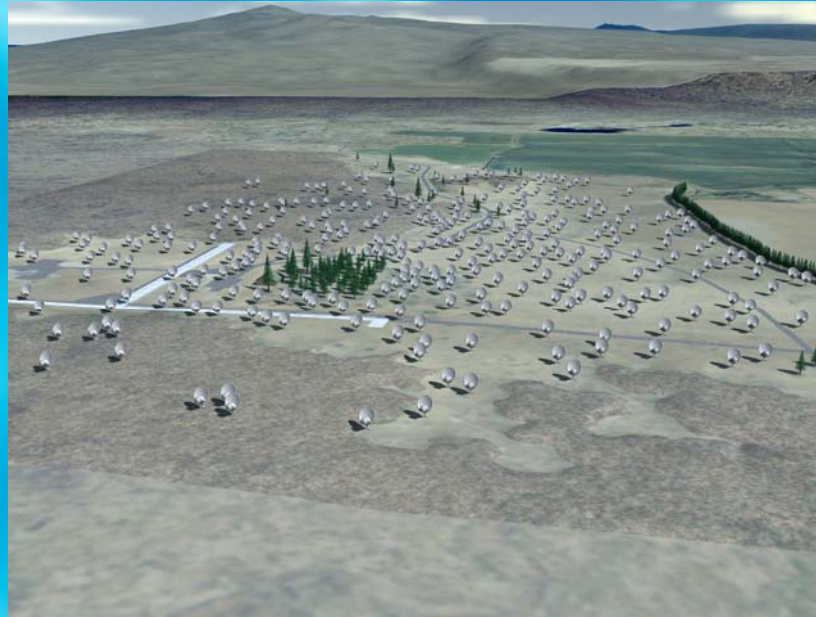
- 1) About 90% of the total flux is from the central region of 2 degree of angular diameter and 50% is within central 50 arcmin region.
- 2) Fluxes increase greatly with the frequency decreases
- 3) Dependency on DM properties and the local magnetic field:

$$S[mJy] \approx 11.25 \langle \sigma v \rangle_{26} B^{4.5} (1 - 0.3B + 0.022B^2)$$

# Observation

- Previous radio continuum observation in the very center region
- No detectable radio emission at  $3\sigma$  noise level of 2mJy were found within 4arcmin of the galaxy centers (Fomalont et al with VLA at 4.885GHz in 1979)

# Allen Telescope Array



- Large field of view (2.5 degrees at 1.4GHz)
  - short baseline
  - 42 working antenna
- ideal for searching diffuse emission from nearby galaxies.

# ATA proposal

Feng Huang, Zhiqiang Shen, Xuelei Chen

Frequency: 1.4 GHz & 750 MHz

Effective Bandwidth: 103 MHz

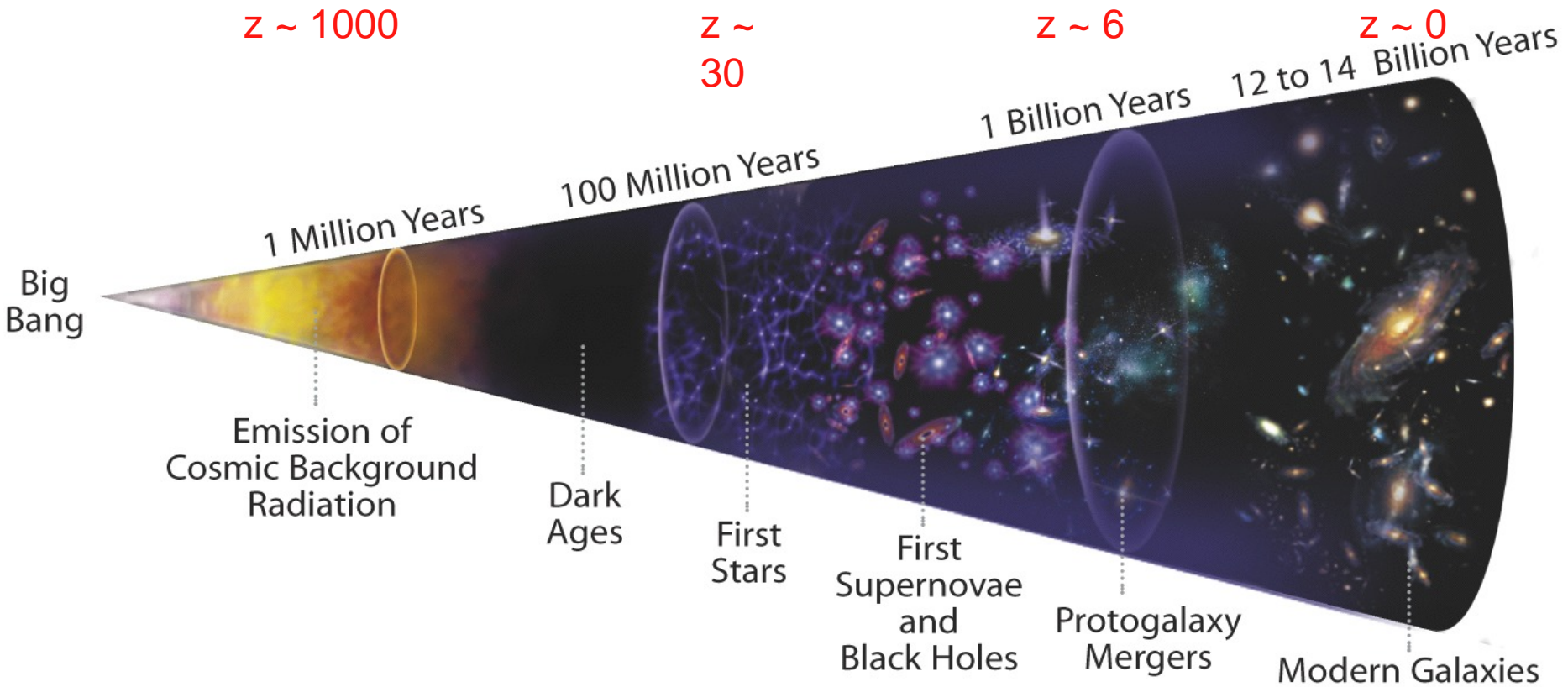
Total Observing Time: 24 hr (half source/half blank field)

Expect: rms 0.05 mJy/beam at 1.42 GHz

Non-detection: will constrain the particle properties of DM and the local environment.

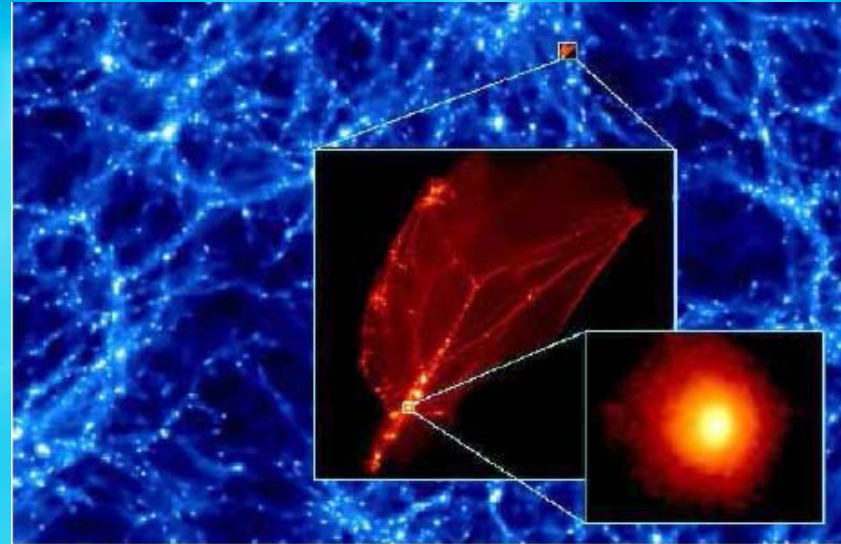
# 21cm signature of first stars

(X.C. & Jordi Miralda-Escude)



# The first stars

- collisionless dark matter collapse to halos
- if gravity exceeds gas pressure gas can fall into halos (Jeans mass)
- gas in the halo can cool
- molecule H cooling  $\sim 10^{2-3}$  K
- atomic cooling  $\sim 10^4$  K



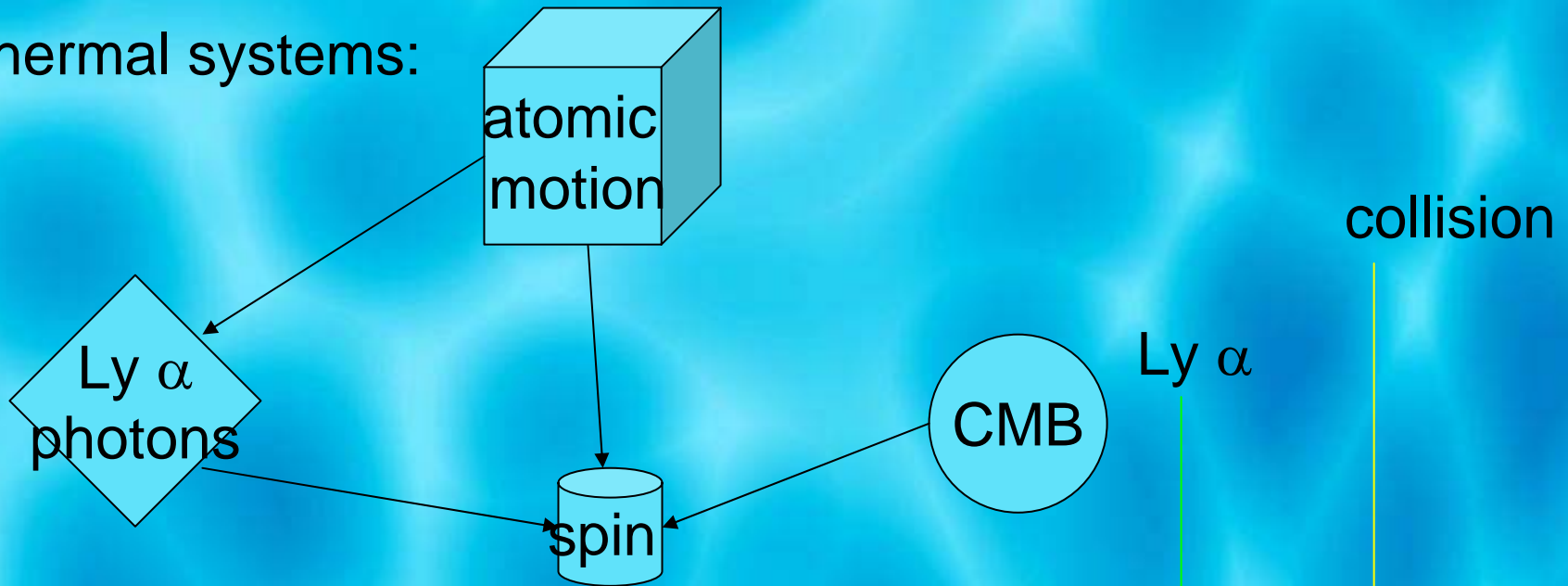
Simulations (Abel 2000, Bromm 2000) indicate first star may form in halos of  $10^{5-6}$  solar mass, one or a few per halo, with masses of a few hundredsolar.

$$M_J \simeq 700M_{\odot} \left( \frac{T}{200 \text{ K}} \right)^{3/2} \left( \frac{n}{10^4 \text{ cm}^{-3}} \right)^{-1/2}$$

# Spin Temperature determines 21cm signature

$$\frac{n_1}{n_0} = 3e^{-\Delta E/k_B T_S} = 3e^{-T_*/T_S}$$

Thermal systems:



$$T_\alpha = T_k$$

$$T_S = \frac{T_{CMB} + y_\alpha T_\alpha + y_c T_k}{1 + y_\alpha + y_c},$$

# The temperature of gas

Heating of the neutral IGM:

- Shock
- UV ionizing radiation (confined to Stromgren sphere)
- Lyman alpha? No  
(Chen & Miralda-Escude 2004, Hirata 2006, chuzhoy & Shapiro 2006, Rybicki 2006, Meiksen 2006, Pritchard & Furlanetto 2006)
- X-ray



# Evolution of global spin temperature

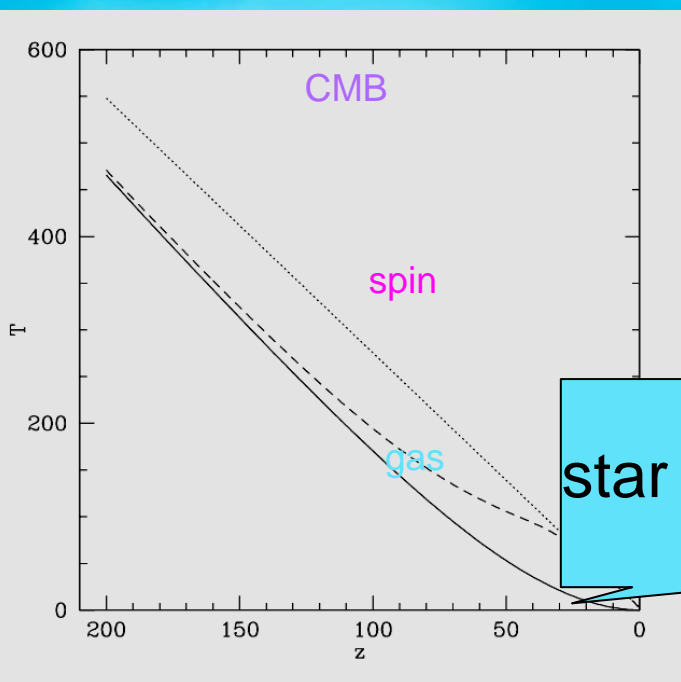
$$z > 150: T_k = T_{\text{cmb}} = T_s$$

$50 < z < 150: T_s = T_k < T_{\text{cmb}}$  collisional coupling

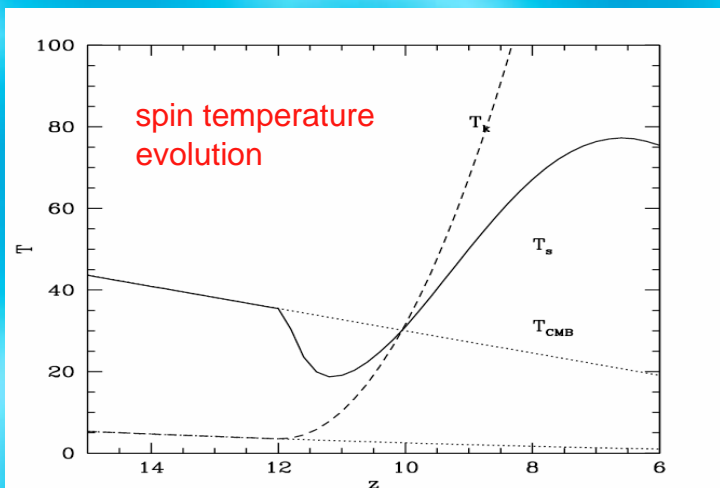
$25 < z < 50: T_k < T_s = T_{\text{cmb}}$  no coupling

$15 < z < 25: T_k < T_s < T_{\text{cmb}}$  Ly alpha coupling

$10 < z < 15: T_k > T_s > T_{\text{cmb}}$  Ly alpha coupling



# Evolution of global spin temperature



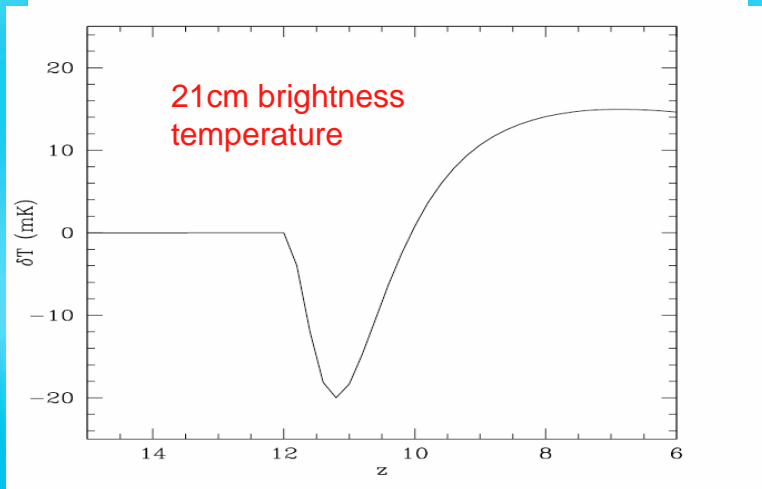
$z > 150$ :  $T_k = T_{\text{cmb}} = T_s$

$50 < z < 150$ :  $T_s = T_k < T_{\text{cmb}}$  collisional coupling

$25 < z < 50$ :  $T_k < T_s = T_{\text{cmb}}$  no coupling

$15 < z < 25$ :  $T_k < T_s < T_{\text{cmb}}$  Ly alpha coupling

$10 < z < 15$ :  $T_k > T_s > T_{\text{cmb}}$  Ly alpha coupling



# Lyman alpha sphere around first stars

first stars: 100 solar mass metal free star radiating at Eddington limit (Bromm et al 2001)

Chen & Miralda-Escude, astro-ph/0605439 (ApJ accepted)

$$T = 1.1 \times 10^5 \left( \frac{M}{100M_{\odot}} \right)^{0.025} K$$
$$L_{total} = 10^{4.5} \frac{M}{M_{\odot}} L_{\odot}$$

- life time of the star  $\sim 3$  Myr, Hubble time  $\sim 10^8$  yr
- light propagation time  $\sim$  size of Ly $\alpha$  sphere (10 kpc)
- halo virial radius  $\sim 0.1$  kpc



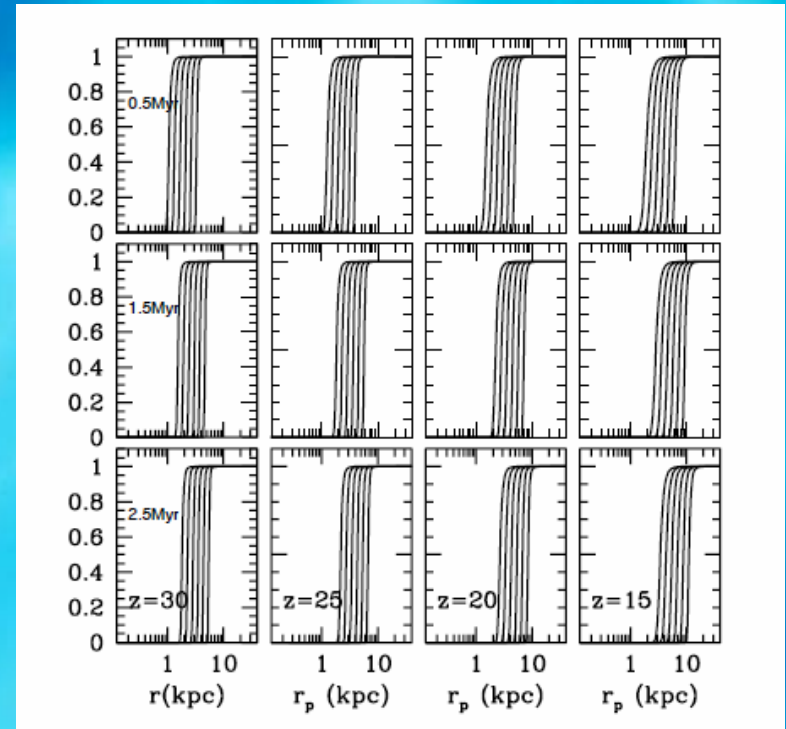
# Radiative Transfer

1D (spherical) radiative transfer code

assume first star blackbody radiation

evolve photon spectrum

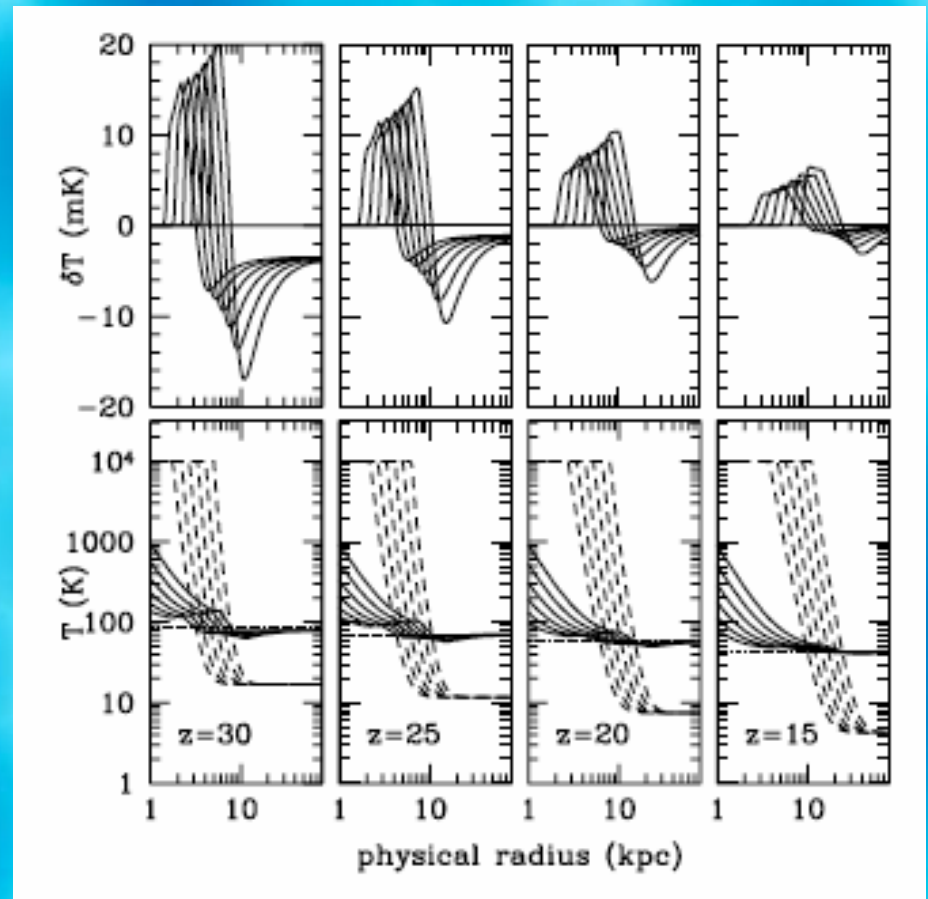
include helium



For stellar mass of 25, 50,  
100, 200, 400, 800  $M_{\text{sun}}$ .

# continuum Lyman alpha photons

- Line photons confined to HII region (or where it is produced) by resonance scattering
- continuum photon: decrease as  $r^{-2}$



# The secondary Lyman alpha photons

induction by X-ray photons:  
recombination, excitation, cascade

photon production rate  $\sim$  energy rate/ $E_\alpha$

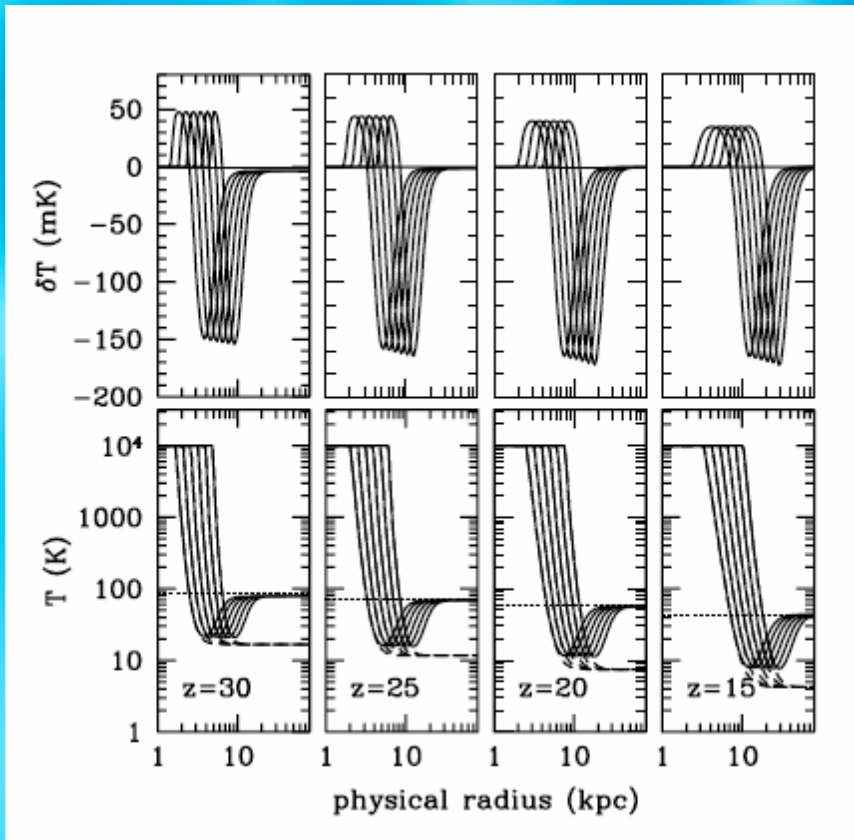
frequency shift rate  $\sim H \nu_\alpha$

flux  $\sim$  photon production rate / frequency shift rate:

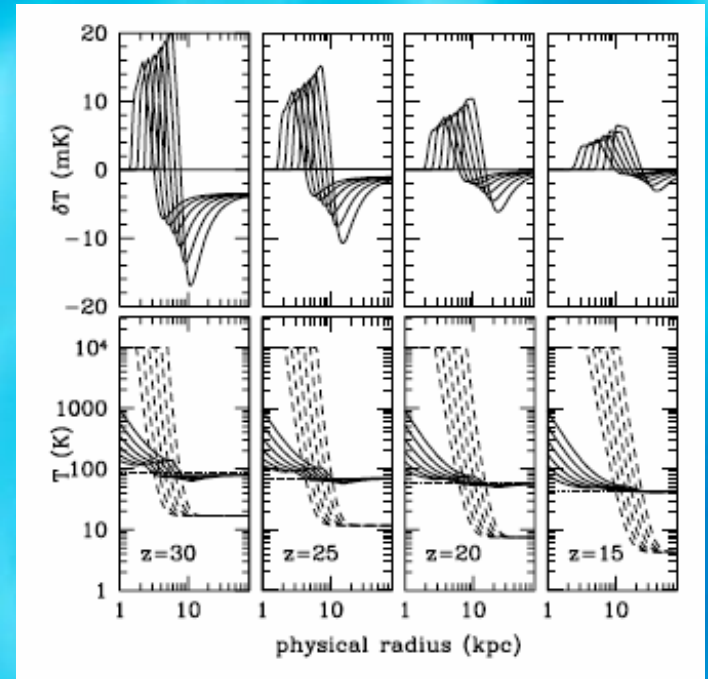
$$J_i = \frac{c\eta_\alpha\Gamma}{4\pi hH\nu_\alpha^2} .$$

Neglected order 1 correction factor and higher Lyman series

# Ly $\alpha$ sphere profile

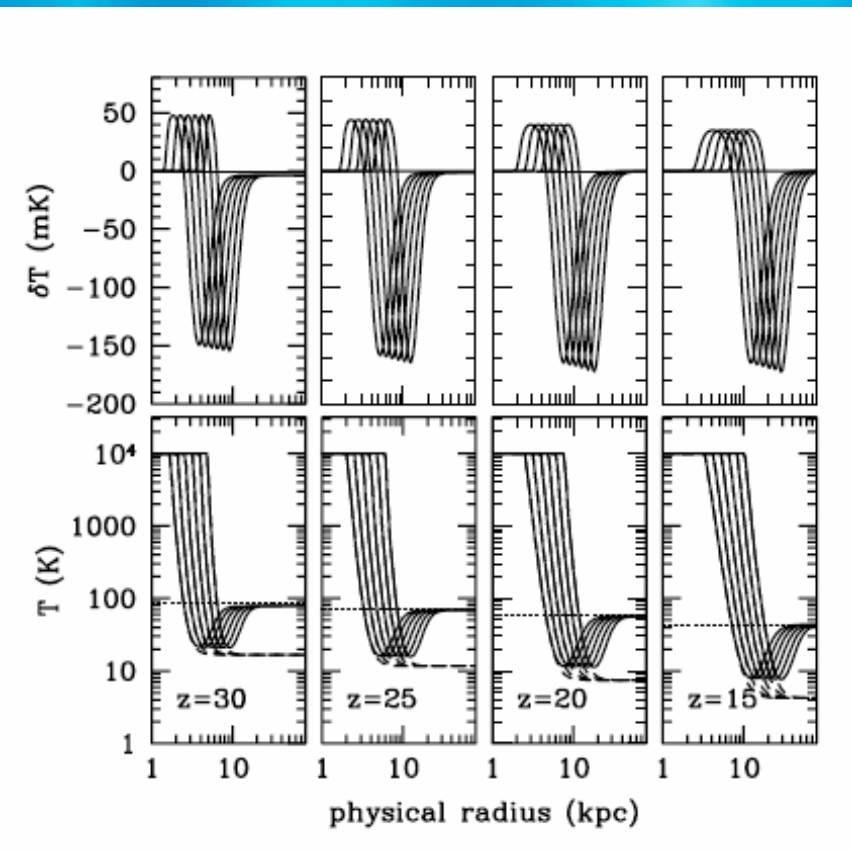


with injected Lyman photons  
very strong absorption

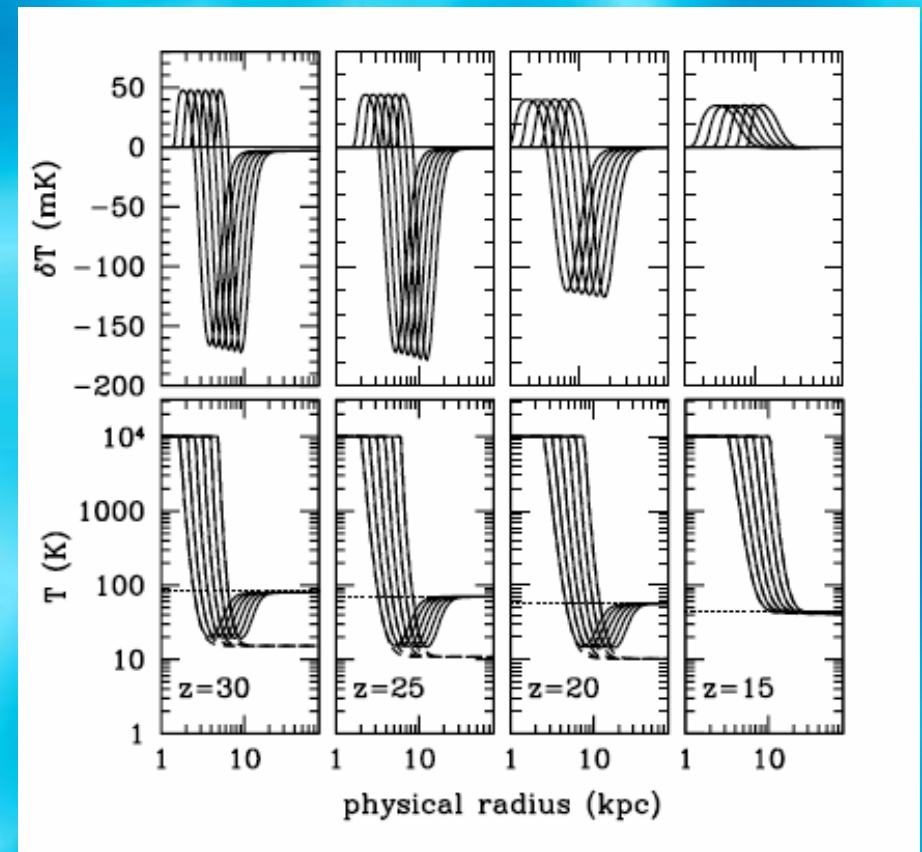


with only continuum Ly  
photons, weak absorption

# The profile of Lyman alpha sphere with X-ray background heating



no heating

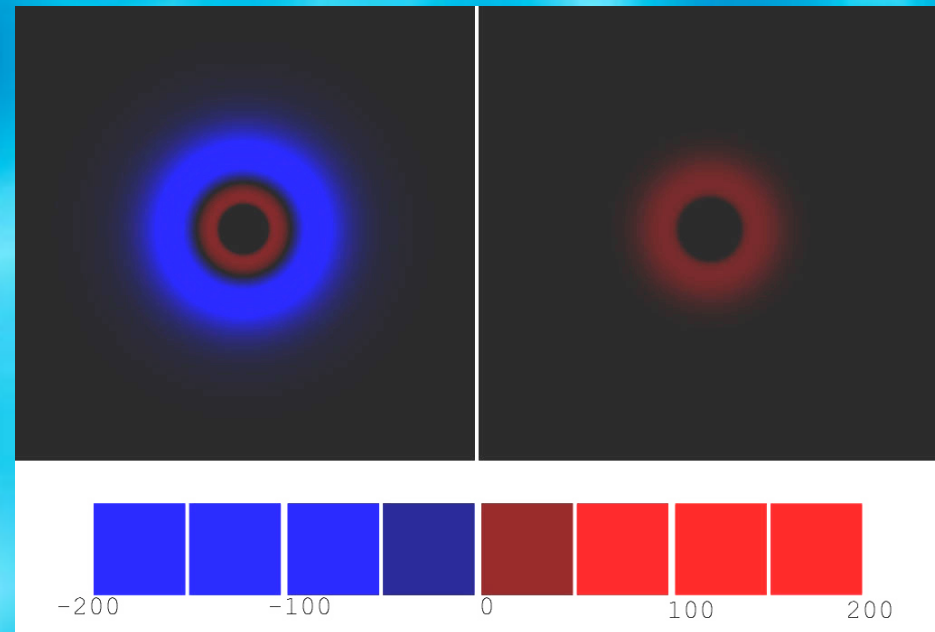


with heating



# Cross Section Map

- Ly alpha background reduce contrast of Ly alpha sphere
- If gas heated above CMB, no absorption signal
- absorption signal much stronger than emission

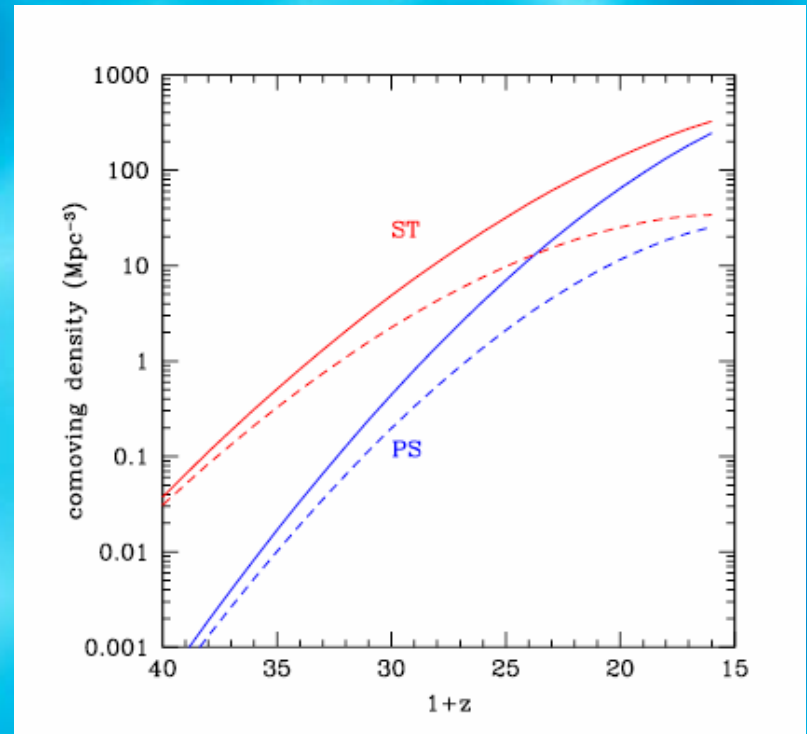


# Formation Rate of first stars

- Minimal mass requirement:  
 $T_{\text{vir}} > 2000 \text{ K}$
- One star formed per halo
- Star died after 3 Myr, so exist only in halos just formed

$$n_* = \int_z^{z_a} dz \int_{M_{\text{min}}(z)}^{\infty} dm \frac{\partial N(m, z)}{\partial z},$$

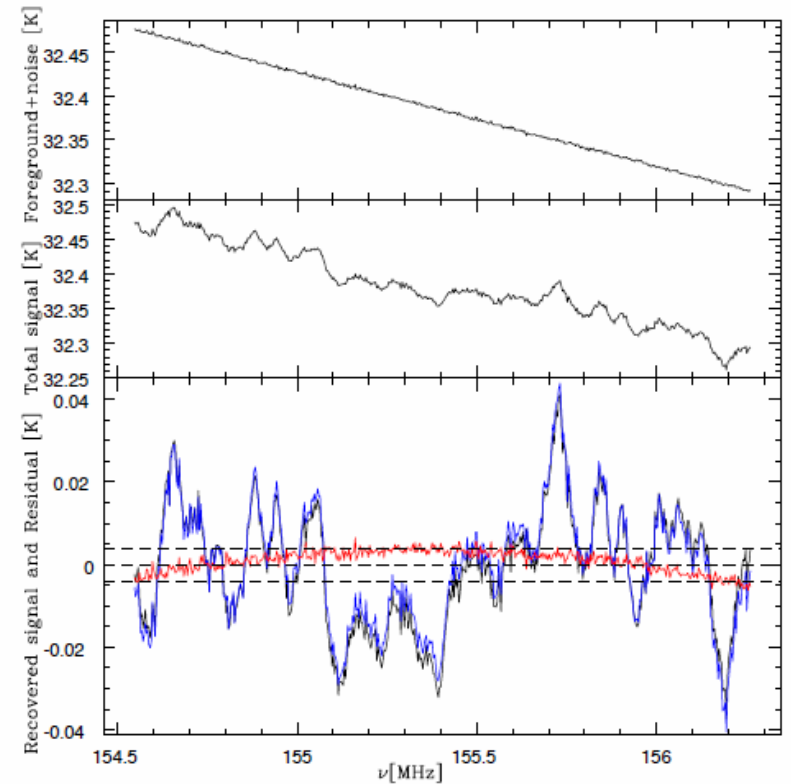
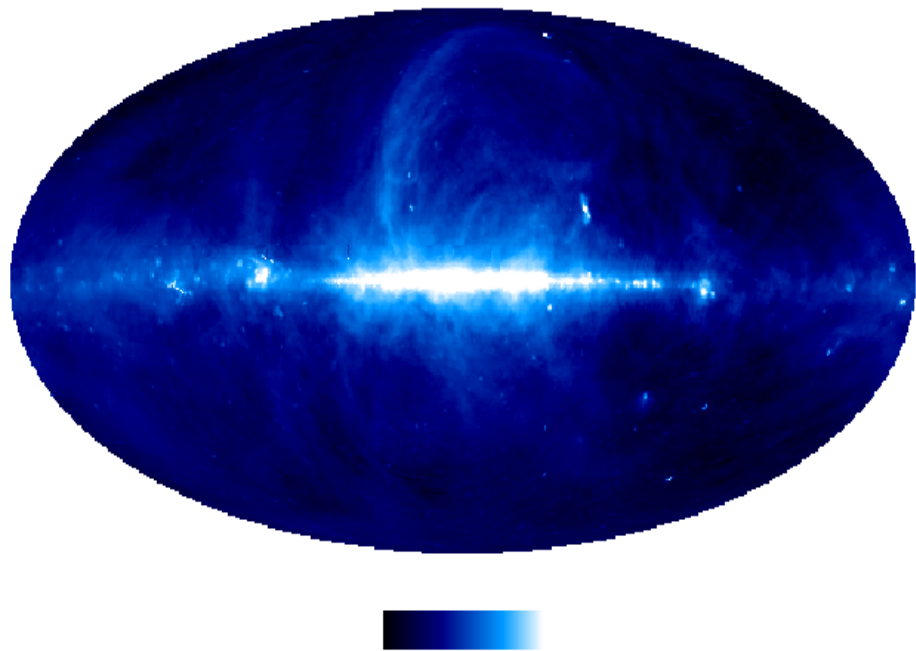
$$n_* \approx (1+z)H(z)t_* \int_{M_{\text{min}}(z)}^{\infty} dm \frac{\partial N(m, z)}{\partial z}$$



This breaks down at low redshift: (1) halo destruction (2) feed back

# Foreground

Galactic Radio Emission



# Observability

difficult to achieve high SNR:  
require almost-filled array

$$\text{SNR} = \sqrt{\Delta\nu t} f_{cov} \frac{\hat{\delta T}}{T_{sys}}$$

$10^5$

1

$10^{-5}$

$$\text{SNR} \sim 3 f_{cov} \left( \frac{1+z}{21} \right)^{-2.5} \left( \frac{\Delta\nu \cdot t}{10 \text{ kHz} \cdot \text{yr}} \right)^{1/2} \left( \frac{\delta T}{10 \text{ mK}} \right)$$

# Some Numbers

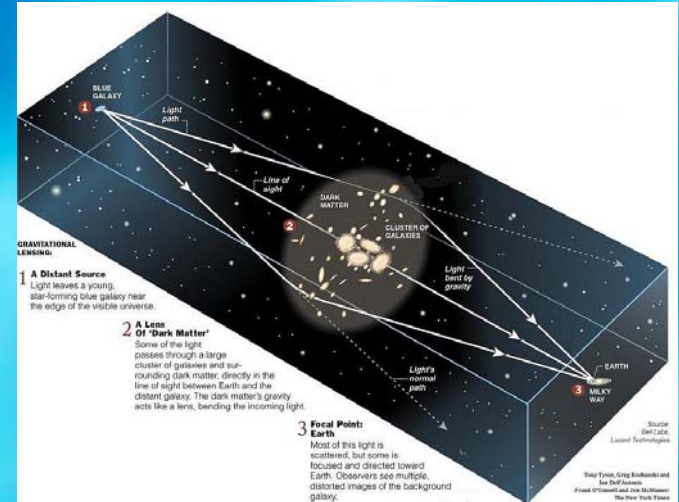
Very Challenging, far beyond the capability of current generation 21cm experiments.

| baseline | bandwidth | beamwidth | SNR |
|----------|-----------|-----------|-----|
| 45 km    | 30 kHz    | 20 arcsec | 5   |
| 65 km    | 30 kHz    | 14 arcsec | 10  |
| 91 km    | 30 kHz    | 10        | 20  |

# Strong gravitational lensing

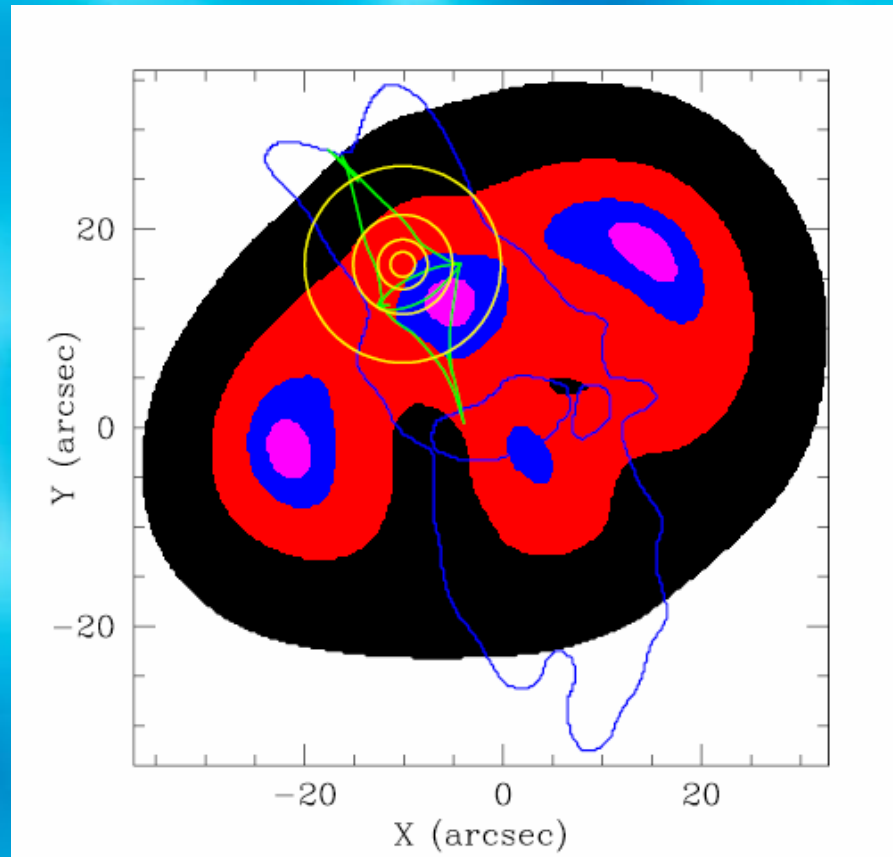
Guoliang Li, Pengjie Zhang, X.C.

- Lensing strong near caustics
- cluster strongly lensed region ~ 10 arcsec size
- significantly increases the size of the image



$$\tau = \frac{1}{4\pi D_s^2} \int_0^{z_s} dz \bar{\sigma}(z) (1+z)^3 \frac{dV_P(z)}{dz},$$

# Can lensing enhance 21cm signal?



G. Li, P. Zhang & XC, astro-ph/0701492, ApJ accepted

# Estimates

Total number of LYA spheres (whole sky)  $10^{11-13}$

optical depth  $10^{-6}$ , a total of  $10^{5-7}$  strongly lensed

The original LYA sphere requires  $A \sim 1000 \text{ km}^2$ , with lensing, reduce to  $A \sim 100 \text{ km}^2$ , still larger than SKA, but there are proposals on constructing such telescopes cheaply (U. Pen & J. Peterson)



# Two aspects in the 21 cm observations

the 21 cm tomography

$$\delta T_b(\nu) \approx 8 x_{HI} (1 + \delta)(1 + z)^{1/2} \left[ 1 - \frac{T_\gamma(z)}{T_s} \right] \left[ \frac{H(z)/(1 + z)}{dv_{\parallel}/dr_{\parallel}} \right] \text{ mK.}$$

the "21 cm forest" (YidongXu, X.C., Zuhui Fan, HyTrac, RenyueCen, work in progress)

$$\tau_\nu(z) = \frac{3h_p c^3 A_{10}}{32\pi k \nu^2} \int_{-\infty}^{+\infty} dr \frac{n_{HI}}{T_s \sqrt{\pi} b(T_k)} \exp \left[ -\frac{u^2(\nu)}{b^2(T_k) + v^2} \right]$$

# The simulation

⊙ reionizationsimulation by (Shin, Trac&Cen 2007)

➤ The ionization slices

- $f_{\text{HI}} \leq 0.5$  - black
- $f_{\text{HI}} > 0.5$  - white

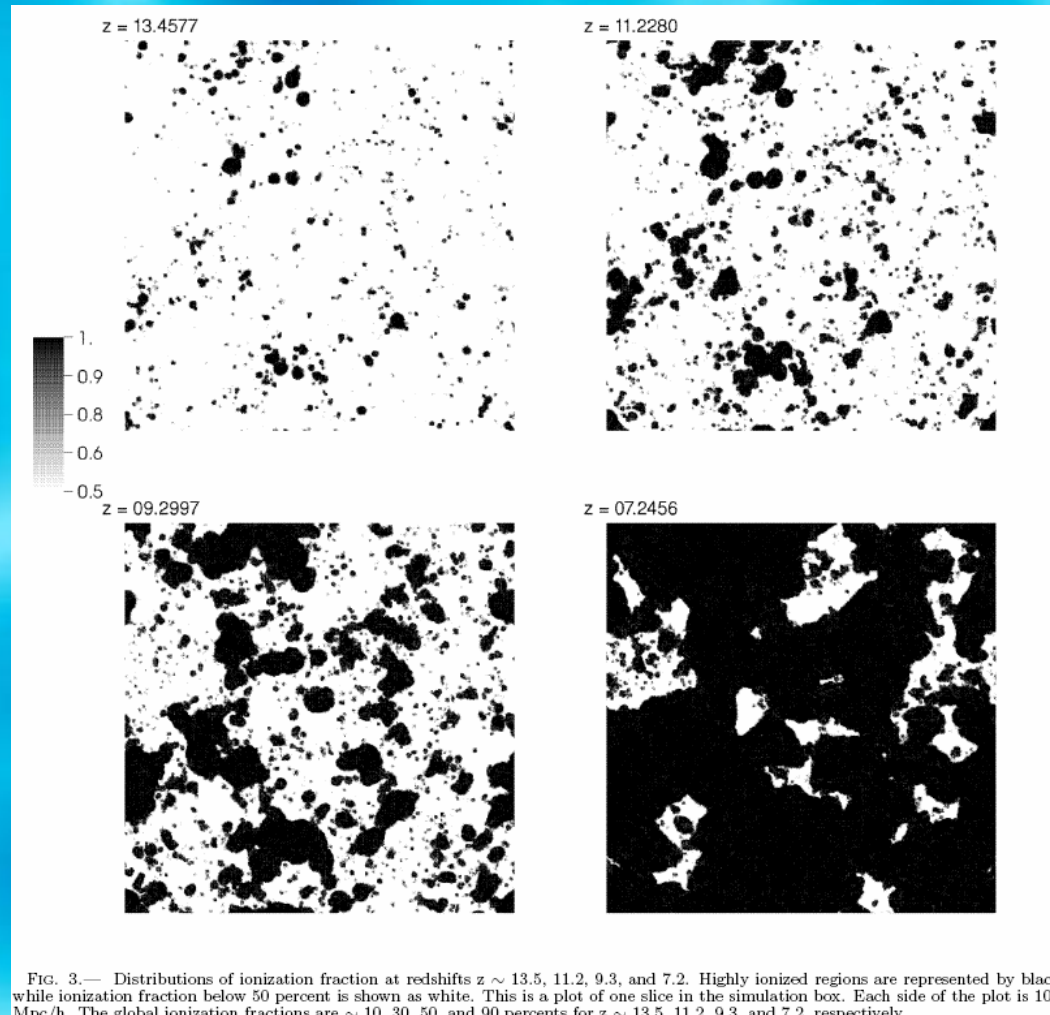
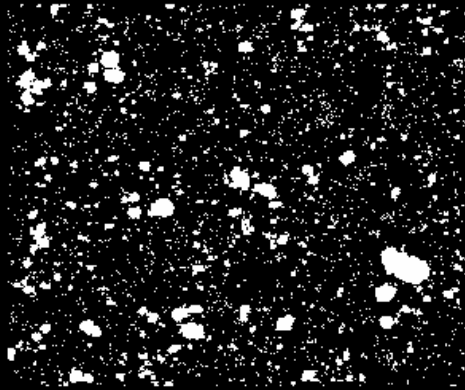


FIG. 3.— Distributions of ionization fraction at redshifts  $z \sim 13.5$ ,  $11.2$ ,  $9.3$ , and  $7.2$ . Highly ionized regions are represented by black while ionization fraction below 50 percent is shown as white. This is a plot of one slice in the simulation box. Each side of the plot is  $10 \text{ Mpc}/h$ . The global ionization fractions are  $\sim 10$ ,  $30$ ,  $50$ , and  $90$  percents for  $z \sim 13.5$ ,  $11.2$ ,  $9.3$ , and  $7.2$ , respectively.

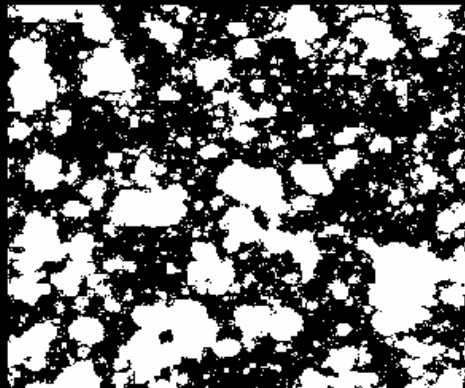
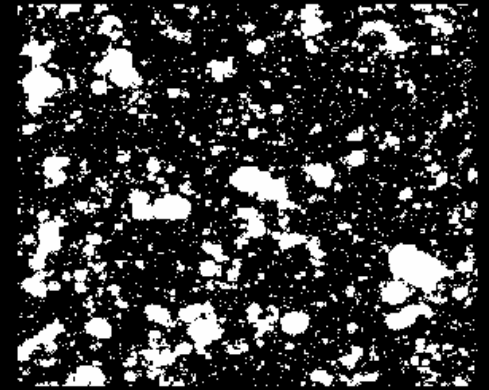
# ◎ The optical depth slices

◆  $\tau$  -slices with relative threshold trace  $\xi$ -slices well, and insensitive to the IGM temperature.

$z = 13.457$



$z = 11.2280$



$z = 09.2997$



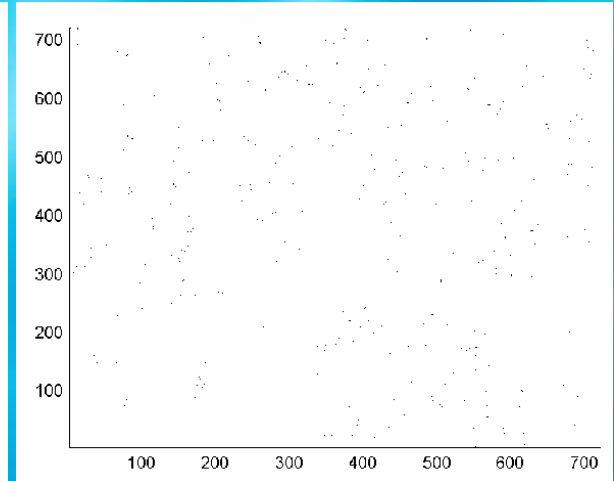
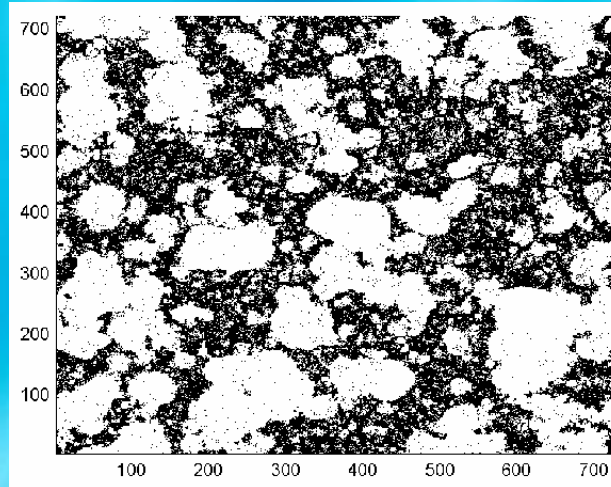
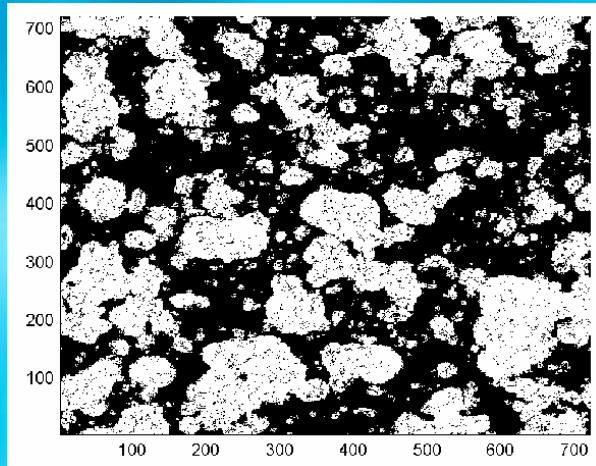
$z = 07.2456$

⊙ The optical depth slices with absolute threshold:

$$T_{\text{IGM}} = 10^1 \text{ K}$$

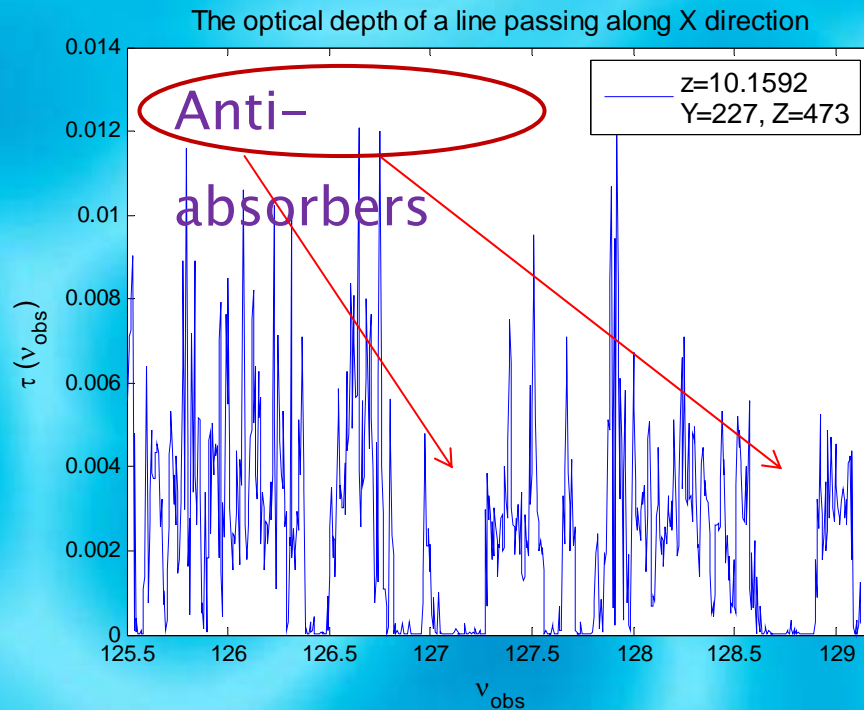
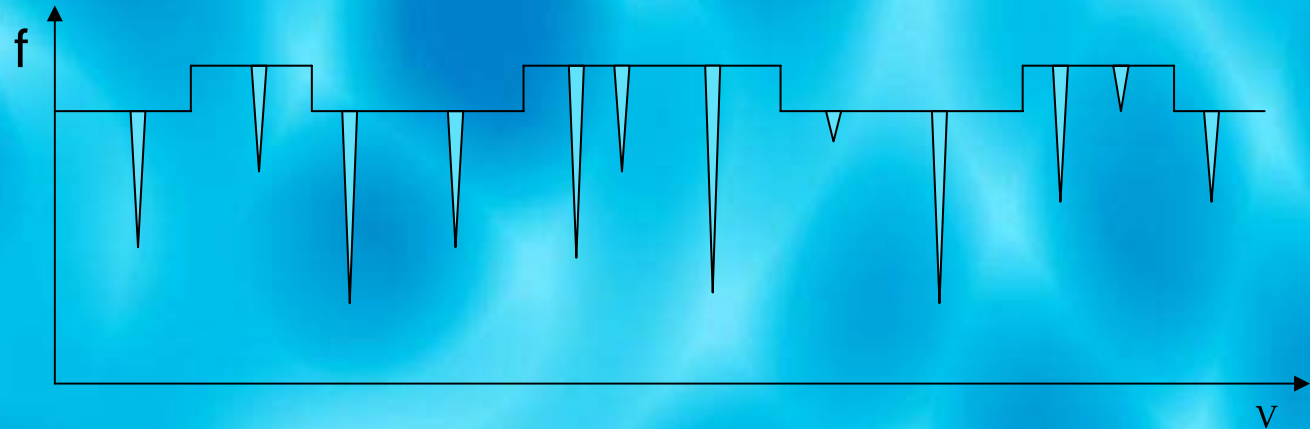
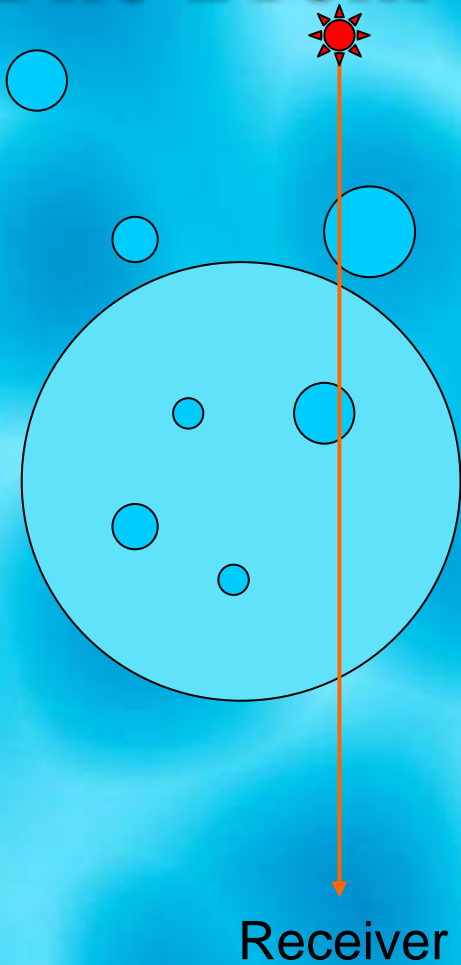
$T_{\text{IGM}} = \text{“real”}$   
temperature from  
simulation ( $\sim 27 \text{ K}$  for  $z =$   
 $8.9$ )

$$T_{\text{IGM}} = 10^2 \text{ K}$$



◆  $\tau$  –slices with absolute threshold are sensitive to the IGM temperature.

# The 21cm Signals

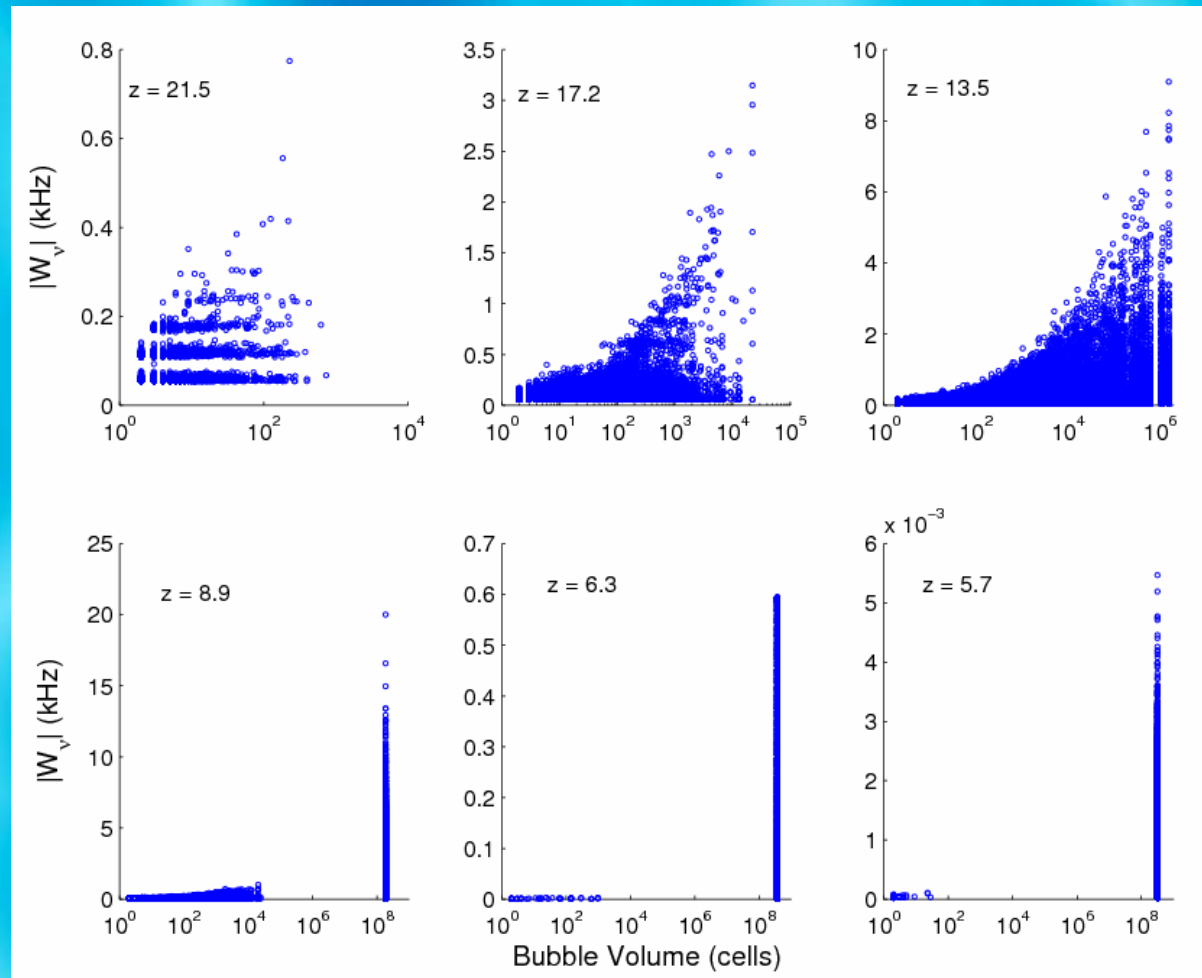


simulation  
data from  
(Shin Trac &  
Cen 2007)

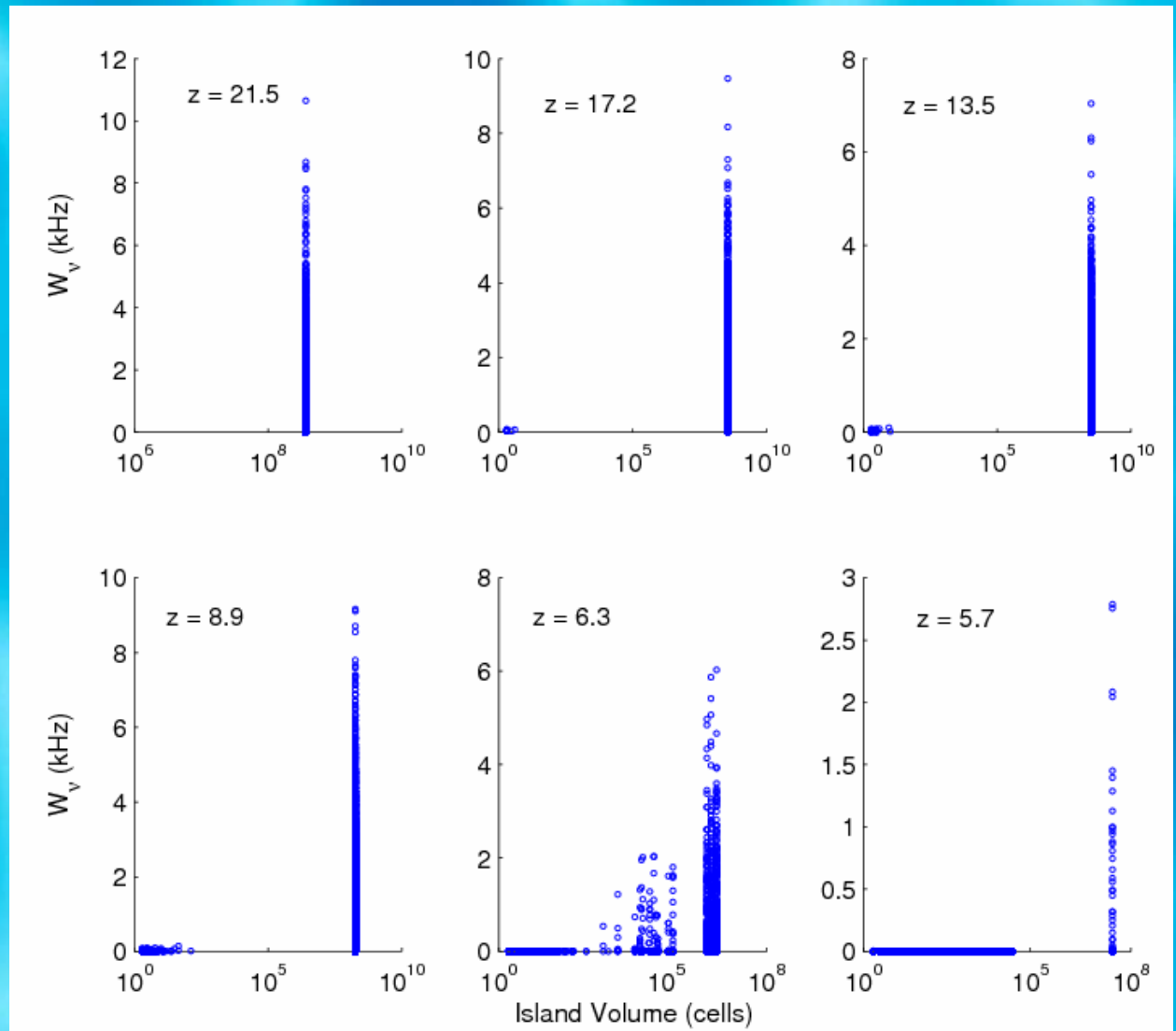
# Equivalent width and bubble volume

- 400 random lines of sight through each side of the simulation box
- Periodic boundary condition is considered
- The threshold:

$$(\tau_0 - \tau_\nu) / \tau_0 \geq 0.5$$



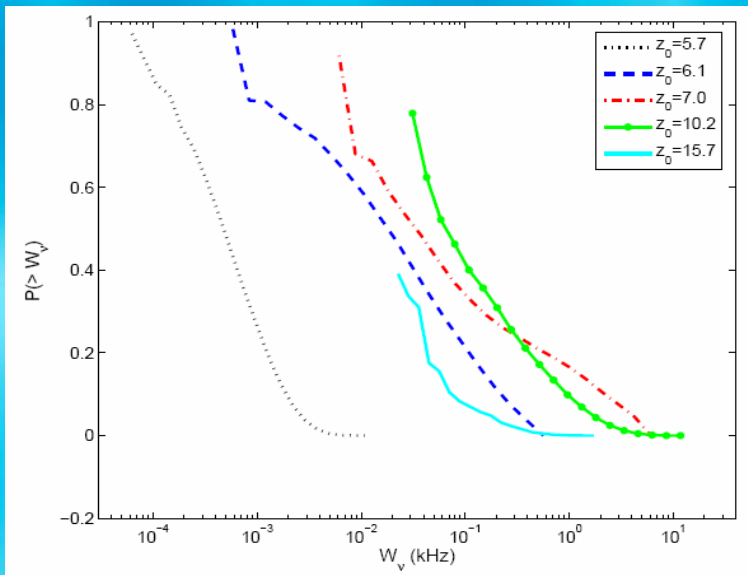
# Equivalent width and island volume



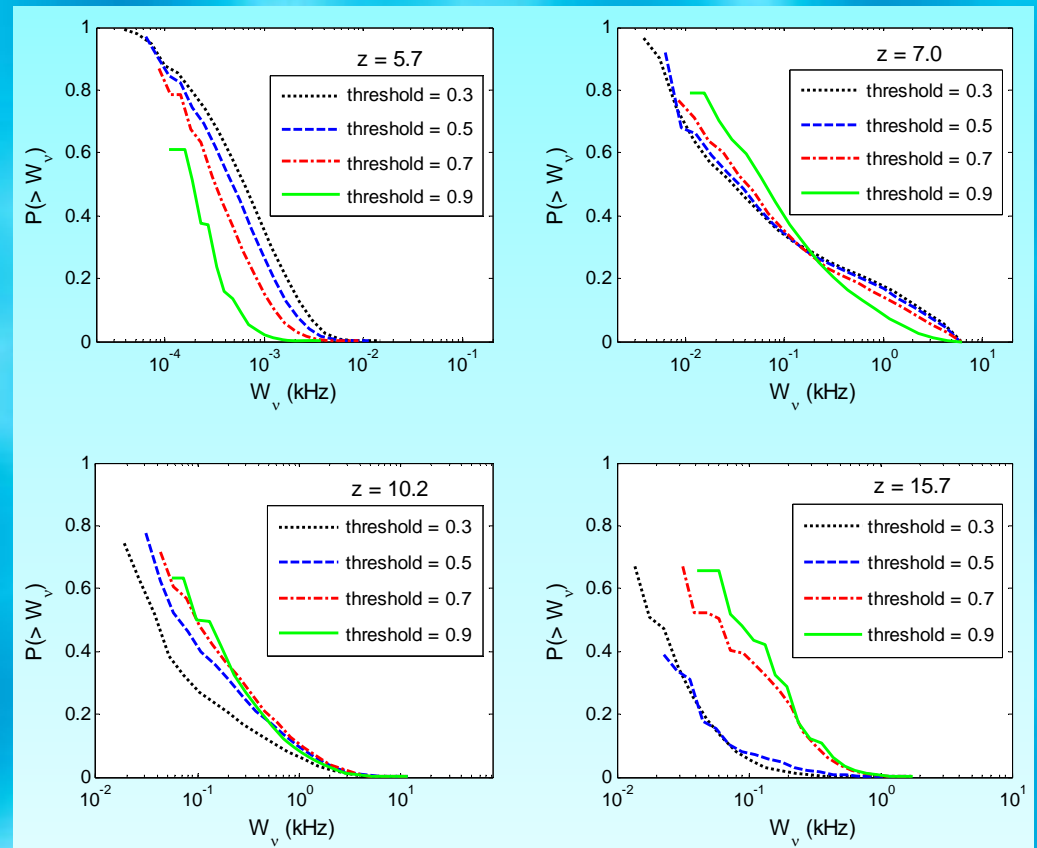
# The “21 cm forest” from simulation

- The cumulative probability distribution (CPD) of the **equivalent widths** for several redshifts.

- ▶ The different thresholds for anti-absorber definition



Show evolution!





Thanks